

# CHEMISTRY IMPROV THEATER

SOMEONE WITH ETHYLENE  
DIBROMIDE RUNS INTO  
SOMEONE WITH PROPYLENE  
DIBROMIDE, AND THE SOLUTIONS  
MIX TOGETHER. HAVE THEM  
DISCUSS THE VAPOR PRESSURE.



# Canonical Equations of Motion -- Hamiltonian Dynamics

- From previous discussion: If the PE  $U$  is independent of velocities (conservative system) then the linear momentum in **rectangular coordinates** has the form:  $\mathbf{p}_i = (\partial L / \partial \dot{\mathbf{x}}_i)$
- Extend this to the case where the Lagrangian is given in terms of generalized coordinates:

**DEFINE: The Generalized Momentum** (corresponding to the generalized coordinate  $q_j$ ):  $\mathbf{p}_j \equiv (\partial L / \partial \dot{q}_j)$ .

$\mathbf{p}_j \equiv$  “Conjugate Momentum” to  $q_j$ .

Or, “Canonically Conjugate Momentum”

- $(q, p) \equiv$  “**conjugate**” or “**canonical**” variables.

- **Lagrange Eqtns of motion:**  $s$  degrees of freedom  

$$(\partial L / \partial \mathbf{q}_j) - (d/dt)[(\partial L / \partial \dot{\mathbf{q}}_j)] = \mathbf{0} \quad (j = 1, 2, 3, \dots, s)$$
- Gives  $s$  **2<sup>nd</sup> order**, time dependent, differential eqtns.  
 $\Rightarrow$  The system motion is determined for all time when **2s initial values** are specified:  $s$   $\mathbf{q}_j$ 's &  $s$   $\dot{\mathbf{q}}_j$ 's
- We can represent the state of the system motion by the time dependent motion of a point in an abstract **s-dimensional configuration space** (coordinates =  $s$  generalized coords  $\mathbf{q}_j$ ).
- **PHYSICS:** In the **Lagrangian Formulation of Mechanics**, a system with  $s$  degrees of freedom = a problem in  $s$  **INDEPENDENT variables**  $\mathbf{q}_j(\mathbf{t})$ . The generalized velocities,  $\dot{\mathbf{q}}_j(\mathbf{t})$  are determined by taking the time derivatives of the  $\mathbf{q}_j(\mathbf{t})$ . The velocities are not independent variables!

# CONTRAST!

- **Hamiltonian Mechanics** is a fundamentally different picture!
  - It describes the system motion in terms of *1<sup>st</sup> order*, time dependent equations of motion. The number of initial conditions is, of course, still **2s**.
- ⇒ We must describe the system motion with **2s independent 1<sup>st</sup> order**, time dependent, differential equations expressed in terms of **2s independent variables**.
- Choose s of these = **s generalized coordinates**  $q_j$ .
  - The others = **s generalized (conjugate) momenta**  $p_j$ .

- **Hamiltonian Mechanics:**
- Describes the system motion in terms of **s** **generalized coordinates  $q_j$**  & **s** **generalized momenta  $p_j$** . Gets 2s *1<sup>st</sup> order*, time dependent equations of motion.
- Recall again the **DEFINITION:** **Generalized Momentum** associated with generalized coord  $q_j$ :
$$p_j \equiv (\partial L / \partial \dot{q}_j)$$
- $(q, p) \equiv$  “conjugate” or “canonical” variables.

# Legendre Transformations

- *Physically*, the Lagrange formulation assumes the coordinates  $\mathbf{q}_j$  are independent variables & the velocities  $\dot{\mathbf{q}}_j$  are dependent variables & are only obtained by taking time derivatives of the  $\mathbf{q}_j$  once the problem is solved.
- *Mathematically*, the Lagrange formalism treats  $\mathbf{q}_j$  &  $\dot{\mathbf{q}}_j$  as independent variables. e.g., in Lagrange's eqtns,  $(\partial L / \partial \mathbf{q}_j)$  means take the partial derivative of  $L$  with respect to  $\mathbf{q}_j$  keeping all other  $\mathbf{q}$ 's & *ALSO* all  $\dot{\mathbf{q}}$ 's constant. Similarly  $(\partial L / \partial \dot{\mathbf{q}}_j)$  means take the partial derivative of  $L$  with respect to  $\dot{\mathbf{q}}_j$  keeping all other  $\dot{\mathbf{q}}$ 's & *ALSO* all  $\mathbf{q}$ 's constant.
- Treated as a pure mathematical problem, changing from the Lagrange formulation to the Hamilton formulation corresponds to changing variables from  $(\mathbf{q}, \dot{\mathbf{q}}, \mathbf{t})$  ( $\mathbf{q}, \dot{\mathbf{q}}$ , independent) to  $(\mathbf{q}, \mathbf{p}, \mathbf{t})$  ( $\mathbf{q}, \mathbf{p}$  independent)

- **Lagrange's Equations** are:

$$(\partial L / \partial \mathbf{q}_j) = (d/dt)[(\partial L / \partial \dot{\mathbf{q}}_j)] \quad (1)$$

- The **generalized momentum** is:  $\mathbf{p}_j = (\partial L / \partial \dot{\mathbf{q}}_j)$  (2)

- (1) & (2) together  $\Rightarrow \dot{\mathbf{p}}_j = (\partial L / \partial \mathbf{q}_j)$  (N's 2<sup>nd</sup> Law!!) (3)

- From previous discussion, the *Hamiltonian* is:

$$\mathbf{H} \equiv \sum_j \dot{\mathbf{q}}_j (\partial L / \partial \dot{\mathbf{q}}_j) - L \quad (4)$$

- Rewrite (4) using the generalized momentum definition (2):

$$\Rightarrow \mathbf{H} = \sum_j \mathbf{p}_j \dot{\mathbf{q}}_j - L \quad (5)$$

- (5) is the starting point for **Hamiltonian Dynamics**.

## Derivation of Hamilton's Equations

- **Lagrangian Dynamics:** Assumes the Lagrangian  $L = L(\mathbf{q}_j, \dot{\mathbf{q}}_j, \mathbf{t})$  (is a function of generalized coordinates, generalized velocities, and time).
- Solve each  $\mathbf{p}_j = (\partial L / \partial \dot{\mathbf{q}}_j)$  for the generalized velocities & get:  $\dot{\mathbf{q}}_j = \dot{\mathbf{q}}_j(\mathbf{q}_k, \mathbf{p}_k, \mathbf{t})$
- In the Hamiltonian  $\mathbf{H} = \sum_j \mathbf{p}_j \dot{\mathbf{q}}_j - L$ , make the change of variables from the set  $(\mathbf{q}_j, \dot{\mathbf{q}}_j, \mathbf{t})$  to the set  $(\mathbf{q}_k, \mathbf{p}_k, \mathbf{t})$  & express the Hamiltonian as:

$$\mathbf{H} = \mathbf{H}(\mathbf{q}_k, \mathbf{p}_k, \mathbf{t}) = \sum_k \mathbf{p}_k \dot{\mathbf{q}}_k - L(\mathbf{q}_k, \dot{\mathbf{q}}_k, \mathbf{t})$$



- **Hamiltonian Dynamics:** Always write

$$\mathbf{H} = \mathbf{H}(\mathbf{q}_k, \mathbf{p}_k, \mathbf{t}) = \sum_k \mathbf{p}_k \dot{\mathbf{q}}_k - L(\mathbf{q}_k, \dot{\mathbf{q}}_k, \mathbf{t})$$

- This stresses that **Hamiltonian dynamics considers the generalized coordinates & generalized momenta as the variables:**  $(\mathbf{q}_k, \mathbf{p}_k, \mathbf{t})$ , whereas **Lagrangian dynamics** considers the generalized coordinates & generalized velocities as the variables:  $(\mathbf{q}_k, \dot{\mathbf{q}}_k, \mathbf{t})$ . A vital & important point!! Please keep it in mind!

- A proper **Hamiltonian** (for use in Hamiltonian dynamics) is **ALWAYS** (!) written as a function of the generalized coordinates & momenta:  $\mathbf{H} \equiv \mathbf{H}(\mathbf{q}_k, \mathbf{p}_k, \mathbf{t})$ . **If you write it as a function of  $\dot{q}_j$ , it is NOT a Hamiltonian!!!**
- A proper **Lagrangian** (for use in Lagrangian dynamics) is **ALWAYS** (!) written as a function of the generalized coordinates & velocities:  $L \equiv L(\mathbf{q}_j, \dot{\mathbf{q}}_j, \mathbf{t})$

## TO EMPHASIZE THIS

Consider a single free particle ( $\vec{p} = m\vec{v}$ )

Energy = **KE** = **T** =  $(1/2)mv^2$  only. So, **H** = **T**

But, if it is a **PROPER HAMILTONIAN**, can  
it be written **H** =  $\frac{1}{2}mv^2$  ?

**NO!!!!!!** **H** *must* be expressed in terms of  
the momentum **p**, not the velocity **v**! So the

**PROPER HAMILTONIAN** is

$$\mathbf{H} = [\mathbf{p}^2/(2m)] \quad \mathbf{!!!!}$$

## Derivation of Hamilton's Equations of Motion

- The Hamiltonian is:  $\mathbf{H} \equiv \mathbf{H}(\mathbf{q}_k, \mathbf{p}_k, \mathbf{t})$

- Take the total differential of  $\mathbf{H}$ :

$$d\mathbf{H} = \sum_k [(\partial\mathbf{H}/\partial\mathbf{q}_k)d\mathbf{q}_k + (\partial\mathbf{H}/\partial\mathbf{p}_k)d\mathbf{p}_k] + (\partial\mathbf{H}/\partial\mathbf{t})d\mathbf{t} \quad (1)$$

- Also:

$$\mathbf{H}(\mathbf{q}_k, \mathbf{p}_k, \mathbf{t}) = \sum_k \mathbf{p}_k \dot{\mathbf{q}}_k - L(\mathbf{q}_k, \dot{\mathbf{q}}_k, \mathbf{t}) \quad (2)$$

- The total differential of (2) is:

$$d\mathbf{H} = \sum_k [\dot{\mathbf{q}}_k d\mathbf{p}_k + \mathbf{p}_k d\dot{\mathbf{q}}_k - (\partial L/\partial\mathbf{q}_k)d\mathbf{q}_k - (\partial L/\partial\dot{\mathbf{q}}_k)d\dot{\mathbf{q}}_k] - (\partial L/\partial\mathbf{t})d\mathbf{t} \quad (3)$$

- We had:  $\mathbf{p}_k = (\partial L / \partial \dot{\mathbf{q}}_k)$  and  $\dot{\mathbf{p}}_k = (\partial L / \partial \mathbf{q}_k)$
- Put these into (3) & get:

$$dH = \sum_k [\dot{\mathbf{q}}_k d\mathbf{p}_k + \mathbf{p}_k d\dot{\mathbf{q}}_k - \mathbf{p}_k d\dot{\mathbf{q}}_k - \dot{\mathbf{p}}_k d\mathbf{q}_k] - (\partial L / \partial t) dt$$

Or:

$$dH = \sum_k [\dot{\mathbf{q}}_k d\mathbf{p}_k - \dot{\mathbf{p}}_k d\mathbf{q}_k] - (\partial L / \partial t) dt \quad (4)$$

- Compare (4) with (1) & get:

$$\dot{\mathbf{q}}_k = (\partial H / \partial \mathbf{p}_k), \quad \dot{\mathbf{p}}_k = - (\partial H / \partial \mathbf{q}_k)$$

$$(\partial H / \partial t) = - (\partial L / \partial t)$$

# Hamilton's Equations of Motion

$$\dot{\mathbf{q}}_k = (\partial\mathbf{H}/\partial\mathbf{p}_k) \qquad \dot{\mathbf{p}}_k = - (\partial\mathbf{H}/\partial\mathbf{q}_k)$$

• We also have:  $(\partial\mathbf{H}/\partial t) = - (\partial L/\partial t)$

• Use **Hamilton's Equations** in (1):

$$d\mathbf{H} = \sum_k [(\partial\mathbf{H}/\partial\mathbf{q}_k)d\mathbf{q}_k + (\partial\mathbf{H}/\partial\mathbf{p}_k)d\mathbf{p}_k] + (\partial\mathbf{H}/\partial t)dt \quad (1)$$

• From Hamilton's Equations, each term in the sum = 0 & (1) becomes:  $(d\mathbf{H}/dt) = (\partial\mathbf{H}/\partial t) \quad (5)$

• (5)  $\Rightarrow$  If there is no *explicit* time dependence in  $\mathbf{H}$ ,  $(\partial\mathbf{H}/\partial t) = \mathbf{0}$ , then total time derivative  $(d\mathbf{H}/dt) = \mathbf{0}$  or  **$\mathbf{H} = \text{constant in time.}$**  ( **$\mathbf{H}$  is conserved.**) If conditions are such that  $\mathbf{H} = \mathbf{E}$  (total energy), then the total energy is conserved!

- *Hamilton's Equations of Motion*

$$\dot{\mathbf{q}}_{\mathbf{k}} = (\partial\mathbf{H}/\partial\mathbf{p}_{\mathbf{k}}) \quad \dot{\mathbf{p}}_{\mathbf{k}} = - (\partial\mathbf{H}/\partial\mathbf{q}_{\mathbf{k}})$$

- These are called **the “canonical” equations of motion** because of their (almost) symmetric appearance.
- The description of motion based on these is called **Hamiltonian Dynamics**.

- *Hamilton's Equations of Motion*

$$\dot{\mathbf{q}}_{\mathbf{k}} = (\partial\mathbf{H}/\partial\mathbf{p}_{\mathbf{k}}) \quad \dot{\mathbf{p}}_{\mathbf{k}} = - (\partial\mathbf{H}/\partial\mathbf{q}_{\mathbf{k}})$$

- Consider a system with  $3\mathbf{n}$  coordinates,  $\mathbf{m}$  constraints &  $s = 3\mathbf{n} - \mathbf{m}$  degrees of freedom.

⇒ **Lagrangian Dynamics** description is in terms of  $s$  Lagrange's Eqtns of Motion (2<sup>nd</sup> order differential eqtns).

⇒ **Hamiltonian Dynamics** description is in terms of  $2s$  Hamilton's Eqtns of Motion (1<sup>st</sup> order differential eqtns).

- Since one can be derived from the other, obviously, the 2 descriptions are **100% equivalent**. They are also 100% equivalent to Newton's 2<sup>nd</sup> Law!



# Discussion of Hamilton's Eqtns

- **Hamiltonian:**  $H(\mathbf{q}, \mathbf{p}, t) = \sum_{\mathbf{k}} \dot{\mathbf{q}}_{\mathbf{k}} \mathbf{p}_{\mathbf{k}} - L(\mathbf{q}, \dot{\mathbf{q}}, t)$  (a)

## Hamilton's Equations of Motion:

$$\dot{\mathbf{q}}_{\mathbf{k}} = (\partial H / \partial \mathbf{p}_{\mathbf{k}}) \quad (\text{b}), \quad \dot{\mathbf{p}}_{\mathbf{k}} = -(\partial H / \partial \mathbf{q}_{\mathbf{k}}) \quad (\text{c}), \quad -(\partial L / \partial t) = (\partial H / \partial t) \quad (\text{d})$$

- **2n 1<sup>st</sup> order**, time dependent equations of motion **replacing** the **n 2<sup>nd</sup> order** Lagrange Equations of motion.
- **(a):** The formal definition of the Hamiltonian **H** in terms of the Lagrangian **L**. However, as we'll see, in practice, we often needn't know **L** first to be able to construct **H**.
- **(b):**  $\dot{\mathbf{q}}_{\mathbf{k}} = (\partial H / \partial \mathbf{p}_{\mathbf{k}})$ : Gives  $\dot{\mathbf{q}}_{\mathbf{k}}$ 's as functions of  $(\mathbf{q}, \mathbf{p}, t)$ .  
 $\Rightarrow$  Given initial values, integrate to get  $\mathbf{q}_{\mathbf{k}} = \mathbf{q}_{\mathbf{k}}(\mathbf{q}, \mathbf{p}, t) \Rightarrow$   
Form the "inverse" of relns of the eqtns  $\mathbf{p}_{\mathbf{k}} = (\partial L / \partial \dot{\mathbf{q}}_{\mathbf{k}})$   
which give  $\mathbf{p}_{\mathbf{k}} = \mathbf{p}_{\mathbf{k}}(\mathbf{q}, \dot{\mathbf{q}}, t)$ .  $\Rightarrow$  "No new information".

- **Hamiltonian:**  $H(\mathbf{q}, \mathbf{p}, t) = \sum_{\mathbf{k}} \dot{\mathbf{q}}_{\mathbf{k}} \mathbf{p}_{\mathbf{k}} - L(\mathbf{q}, \dot{\mathbf{q}}, t)$  (a)

### Hamilton's Equations of Motion:

$$\dot{\mathbf{q}}_{\mathbf{k}} = (\partial H / \partial \mathbf{p}_{\mathbf{k}}) \quad (\text{b}), \quad \dot{\mathbf{p}}_{\mathbf{k}} = -(\partial H / \partial \mathbf{q}_{\mathbf{k}}) \quad (\text{c}), \quad -(\partial L / \partial t) = (\partial H / \partial t) \quad (\text{d})$$

- **(b):**  $\dot{\mathbf{q}}_{\mathbf{k}} = (\partial H / \partial \mathbf{p}_{\mathbf{k}}): \Rightarrow \dot{\mathbf{q}}_{\mathbf{k}} = \dot{\mathbf{q}}_{\mathbf{k}}(\mathbf{q}, \mathbf{p}, t). \Rightarrow$  “No new info”.
  - Usually true in terms of **SOLVING** mechanics problems. However, within the Hamiltonian picture of mechanics, where  $\mathbf{H} = \mathbf{H}(\mathbf{q}, \mathbf{p}, t)$ , obtained **NO MATTER HOW** (not necessarily by (a)), this has equal footing (& contains equally important information as (c)).
- **(c):**  $\dot{\mathbf{p}}_{\mathbf{k}} = -(\partial H / \partial \mathbf{q}_{\mathbf{k}}):$  Integrate & get  $\mathbf{p}_{\mathbf{k}} = \mathbf{p}_{\mathbf{k}}(\mathbf{q}, \mathbf{p}, t)$
- **(d):**  $-(\partial L / \partial t) = (\partial H / \partial t):$  This is obviously only important in time dependent problems.

# Recipe for Hamiltonian Mechanics

- **Hamiltonian:**  $\mathbf{H}(\mathbf{q},\mathbf{p},t) = \sum_{\mathbf{k}} \dot{\mathbf{q}}_{\mathbf{k}} \mathbf{p}_{\mathbf{k}} - L(\mathbf{q},\dot{\mathbf{q}},t)$  (a)

## Hamilton's Equations of Motion:

$$\dot{\mathbf{q}}_i = (\partial\mathbf{H}/\partial\mathbf{p}_i) \text{ (b)}, \dot{\mathbf{p}}_i = -(\partial\mathbf{H}/\partial\mathbf{q}_i) \text{ (c)}, -(\partial L/\partial t) = (\partial\mathbf{H}/\partial t) \text{ (d)}$$

- **Recipe:**

1. Set up the Lagrangian,  $L = T - U = L(\mathbf{q},\dot{\mathbf{q}},t)$
2. Compute s conjugate momenta using:  $\mathbf{p}_{\mathbf{k}} \equiv (\partial L/\partial \dot{\mathbf{q}}_{\mathbf{k}})$
3. Form the Hamiltonian  $\mathbf{H}$  from (a).

This will be of the “mixed” form  $\mathbf{H} = \mathbf{H}(\mathbf{q},\dot{\mathbf{q}},\mathbf{p},t)$

4. Invert s  $\mathbf{p}_{\mathbf{k}} \equiv (\partial L/\partial \dot{\mathbf{q}}_{\mathbf{k}})$  to get  $\dot{\mathbf{q}}_{\mathbf{k}} \equiv \dot{\mathbf{q}}_{\mathbf{k}}(\mathbf{q},\mathbf{p},t)$ .
5. Apply the results of 4. to eliminate  $\mathbf{q}_{\mathbf{k}}$  from  $\mathbf{H}$  to get a ***proper Hamiltonian***  $\mathbf{H} = \mathbf{H}(\mathbf{q},\mathbf{p},t)$ .

***Then & only then*** can you properly & correctly use (b) & (c) to get the equations of motion!

- If you think that this is a long, tedious process, you aren't alone!
  - This requires that you set up the Lagrangian first!
  - If you already have the Lagrangian, why not go ahead & do Lagrangian dynamics instead of going through all of this to do Hamiltonian dynamics?
  - Further, combining the **2s** 1<sup>st</sup> order differential equations of motion

$$\dot{\mathbf{q}}_{\mathbf{k}} = (\partial\mathbf{H}/\partial\mathbf{p}_{\mathbf{k}}) \quad (\mathbf{b}) \quad \& \quad \dot{\mathbf{p}}_{\mathbf{k}} = - (\partial\mathbf{H}/\partial\mathbf{q}_{\mathbf{k}}) \quad (\mathbf{c})$$

gives the **SAME** s 2<sup>nd</sup> order differential equations of motion that Lagrangian dynamics gives!

- However, fortunately, for many physical systems of interest, it's possible to considerably shorten this procedure, even eliminating many steps completely!

## Hamiltonian Dynamics Recipe (Comments)

1. In simple cases, its often possible to directly set up  $\mathbf{H}$  in terms of generalized coordinates & momenta:  $\mathbf{H}(\mathbf{q}_k, \mathbf{p}_k, \mathbf{t})$ . This is because the identification of the generalized momenta  $\mathbf{p}_k$  is easy.
2. Unfortunately, in more complicated cases, where identification of the  $\mathbf{p}_k$  is more difficult, it actually might be necessary to first set up the Lagrangian  $L$ , & then calculate  $\mathbf{p}_k$  with  $\mathbf{p}_k = (\partial L / \partial \dot{\mathbf{q}}_k)$
3. Only after you have  $\mathbf{H} = \mathbf{H}(\mathbf{q}_k, \mathbf{p}_k, \mathbf{t})$  can you apply Hamilton's Eqtns!  $\dot{\mathbf{q}}_k = (\partial \mathbf{H} / \partial \mathbf{p}_k)$ ,  $\dot{\mathbf{p}}_k = - (\partial \mathbf{H} / \partial \mathbf{q}_k)$

- *In many cases*, we know that

$$\Rightarrow \mathbf{H} = \mathbf{T} + \mathbf{U}$$

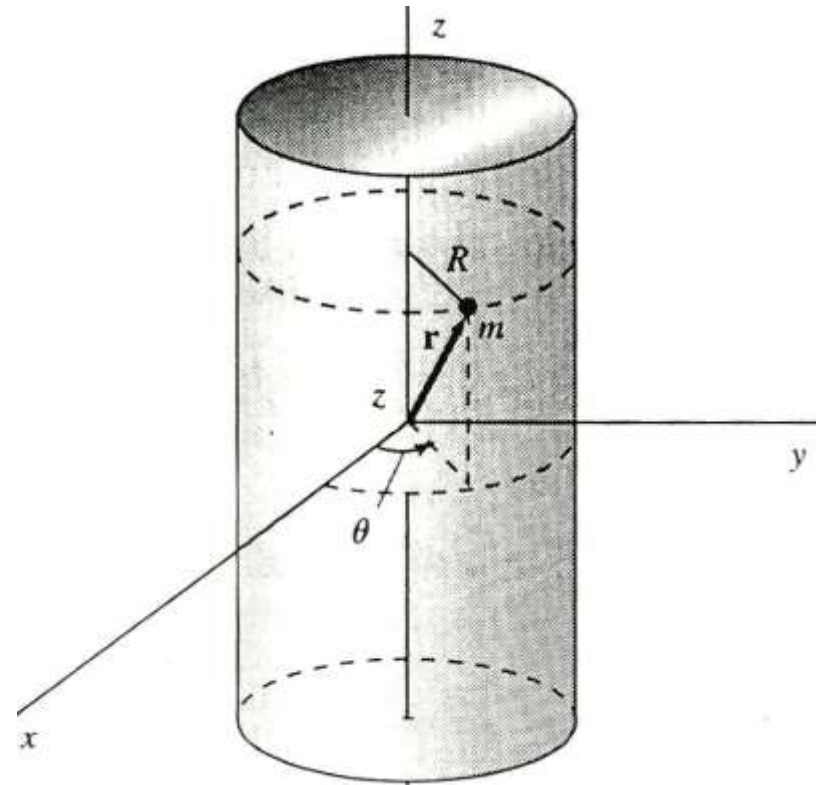
That is, in many cases, the Hamiltonian is automatically the total mechanical energy  $E$

- If that is the case, we can *skip many steps* of the recipe. Instead: 1) Write  $\mathbf{H} = \mathbf{T} + \mathbf{U}$  *immediately*. Express  $\mathbf{T}$  in terms of the MOMENTA  $\mathbf{p}_k$  (not the velocities  $\dot{\mathbf{q}}_k$ ). Often it's easy to see how the  $\mathbf{p}_k$  depend on the  $\dot{\mathbf{q}}_k$  & thus its easy to do this. Once this is done, we can **go ahead & do Hamiltonian Dynamics** without ever having written the Lagrangian.

## Example 7.11

- Use the Hamiltonian method to find the equations of motion of a particle of mass  $m$ , constrained to move on the surface of a cylinder defined by  $x^2 + y^2 = R^2$ . The particle is subject to a force directed towards the origin and proportional to the distance of the particle from the origin:

$$\vec{F} = -k\vec{r}. \text{ Worked on the board!}$$



# Comments

- We could have done this problem with **Lagrangian** mechanics! In fact, in this case, **Lagrange's Equations** are easier to get than **Hamilton's Equations**.
- **This is often true!** The **Lagrangian Method** often leads more easily to the equations of motion than the **Hamiltonian Method**. (In my opinion, almost always!)
- In the **Hamiltonian method**,  $q_k$  &  $p_k$  are considered independent. In the **Lagrangian Method**  $q_k$ ,  $\dot{q}_k$  are not.  
⇒ **The Hamiltonian Method** sometimes has a practical advantage over the **Lagrangian Method**.
  - e. g., Useful in celestial mechanics: Looking at motions of bodies due to perturbations of other bodies.
- In general, main power of **Hamiltonian** formulation is its use as a basis for sub-areas of physics beyond classical mechanics.
  - Quantum mechanics and beyond!